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Grant N00014-90-J-1193

TECHNICAL REPORT No. 18

Specific Heat of Anisotropic Superconductors

by

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Prepared for Publication

in

Quarterly Report of the New York Institute on Superconductivity
Volume 3, Number 1, Summer 1990

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State University of New York at Buffalo
Buffalo, New York 14260

July 1990

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REPORT DOCUMENTATION PAGE

Form Approved
OMB No. 0704-0188

1a. REPORT SECURITY CLASSIFICATION Unclassified			1b. RESTRICTIVE MARKINGS		
2a. SECURITY CLASSIFICATION AUTHORITY			3. DISTRIBUTION/AVAILABILITY OF REPORT Approved for public release; distribution unlimited		
2b. DECLASSIFICATION/DOWNGRADING SCHEDULE					
4. PERFORMING ORGANIZATION REPORT NUMBER(S) UBUFFALO/DC/90/TR-18			5. MONITORING ORGANIZATION REPORT NUMBER(S)		
6a. NAME OF PERFORMING ORGANIZATION Depts. Chemistry & Physics State University of New York		6b. OFFICE SYMBOL (If applicable)		7a. NAME OF MONITORING ORGANIZATION	
6c. ADDRESS (City, State, and ZIP Code) Fronczak Hall, Amherst Campus Buffalo, New York 14260			7b. ADDRESS (City, State, and ZIP Code) Chemistry Program 800 N. Quincy Street Arlington, Virginia 22217		
8a. NAME OF FUNDING/SPONSORING ORGANIZATION Office of Naval Research		8b. OFFICE SYMBOL (If applicable)		9. PROCUREMENT INSTRUMENT IDENTIFICATION NUMBER Grant N00014-90-J-1193	
8c. ADDRESS (City, State, and ZIP Code) Chemistry Program 800 N. Quincy Street Arlington, Virginia 22217			10. SOURCE OF FUNDING NUMBERS		
			PROGRAM ELEMENT NO.	PROJECT NO.	TASK NO.
			WORK UNIT ACCESSION NO.		
11. TITLE (Include Security Classification) Specific Heat of Anisotropic Superconductors					
12. PERSONAL AUTHOR(S) D. Sahu, A. Langner and Thoma, George					
13a. TYPE OF REPORT		13b. TIME COVERED FROM _____ TO _____		14. DATE OF REPORT (Year, Month, Day) July 1990	
				15. PAGE COUNT 5	
16. SUPPLEMENTARY NOTATION Quarterly Report of the New York State Institute on Superconductivity Volume 3, Number 1, Summer 1990					
17. COSATI CODES			18. SUBJECT TERMS (Continue on reverse if necessary and identify by block number)		
FIELD	GROUP	SUB-GROUP	ANISOTROPIC SUPERCONDUCTORS, FERMION SURFACE,		
			SPECIFIC HEAT, DEBYE ENERGY		
			ENERGY GAP, MATSUBARA FREQUENCY		
19. ABSTRACT (Continue on reverse if necessary and identify by block number) A theoretical approach for obtaining a Ginzburg-Landau free-energy expansion for a class of even-parity superconducting states that enables the calculation of thermodynamic quantities for heavy-fermion and other conventional superconductors is discussed.					
20. DISTRIBUTION/AVAILABILITY OF ABSTRACT <input checked="" type="checkbox"/> UNCLASSIFIED/UNLIMITED <input checked="" type="checkbox"/> SAME AS RPT. <input type="checkbox"/> DTIC USERS			21. ABSTRACT SECURITY CLASSIFICATION Unclassified		
22a. NAME OF RESPONSIBLE INDIVIDUAL Dr. David L. Nelson			22b. TELEPHONE (Include Area Code) (202) 696-4410		22c. OFFICE SYMBOL

SPECIFIC HEAT OF ANISOTROPIC SUPERCONDUCTORS

by

D. Sahu, A. Langner and Thomas F. George

Thermodynamic properties of superconductors provide valuable information about the nature of the superconducting state. The existence of an energy gap Δ in isotropic superconductors leads to an exponential dependence of specific heat on the gap function: $C_s \propto \exp(-\Delta/k_B T)$, where k_B is Boltzmann's constant and T is the absolute temperature. For anisotropic superconductors, the energy gap function Δ vanishes along lines or planes of the Fermi surface leading to a power law dependence of specific heat on the temperature. The temperature dependence of specific heat thus leads one to draw conclusions about the nature of the excitation spectrum.

It has been realized in recent years that in "heavy-fermion" superconductors,¹ electron pairing could arise due to a p-wave state or a d-wave state. For a p-wave state, the relative orbital angular momentum of the pair, l , is equal to 1 and the spins form a triplet state. On the other hand, for a d-wave state the relative orbital angular momentum of the pair is given by $l = 2$ and the spins form a singlet state. Calculation of the specific heat for these anisotropic superconductors can be done analogous to Bardeen-Cooper-Schrieffer (BCS) superconductors from the free energy expansion. In the imaginary time representation, the free energy difference, $F_n - F_s$, between the normal state and the superconducting state for a BCS superconductor is²

$$F_n - F_s = N(0) \left\{ -\Delta^2 \ln(t) + 4\pi T \sum_{n=0}^{n_D} \left[(|\Delta|^2 + \omega_n^2)^{1/2} - \omega_n - |\Delta|^2 (2\omega_n)^{-1} \right] \right\},$$

where $N(0)$ is the density of states per single spin evaluated at the Fermi surface, t is the reduced temperature, $\omega_n = \pi T(2n+1)$ is the Matsubara frequency, and n is an integer with an upper limit of n_D corresponding to the Debye cutoff frequency. The entropy S and the specific heat C_s are related to the free energy through $S = - (dF/dT)$ and $C_s = T(dS/dT)$. In carrying out the differentiation, one has to bear in mind that both Δ and ω_n depend on temperature.

The BCS theory is based on a simple intuitive model for the effective attractive interaction between electrons, namely that the interaction potential is attractive for wave vectors k that lie within a small spread of $\pm \Delta_k$ around the Fermi wave vector k_F and that the potential is ≥ 0 otherwise. In addition, the Debye energy ω_D is assumed to be much larger than the energy ξ_k of the electrons. This approximation is called the "weak-coupling" limit, and many superconductors fall into this category. However, there are superconductors for which the electron-phonon coupling is strong, requiring the BCS theory to be modified. This is done in the "strong-coupling" theory of Eliashberg, which utilizes the fact that the electron-phonon interaction is not instantaneous, but rather is delayed in time (retardation). For details of this theory we refer the reader to the book by Mahan.³ Other applications and extensions of the BCS ideas can be found in the excellent two-volume set edited by R. D. Parks.⁴ It should be mentioned here that a strong-coupling, phonon-mediated model has been proposed to explain the properties of the new "high- T_c " superconductors.⁵ However, this model deviates from BCS theory in

that the primary charge carriers are bosons (integral spin) rather than fermions.

The present authors have developed a systematic analysis to obtain a Ginzburg-Landau free-energy expansion for a class of even-parity superconducting states^{6,7} that enables the calculations of thermodynamic quantities for the heavy-fermion and other conventional superconductors. A preliminary application of this approach to model the effect of anisotropy in the new "high-temperature" superconductors has also been pursued with some measure of success.⁸ The main ideas in the above works are: (1) the superconducting states can be anisotropic consistent with the symmetry of the underlying lattice and (2) the anisotropic states can couple to other anisotropic or isotropic states that are permitted by group theory. The construction of these generalized free energies follows the usual prescription of combined operations of rotational invariance, time reversal invariance and gauge invariance. The Matsubara imaginary time representation then provides a convenient tool for carrying out the analytical calculation of the thermodynamic quantities. This tool is again very convenient to handle complex situations, such as evaluation of position-dependent energy gaps that occur when a bulk superconductor is in contact with a thin film of a normal metal (a proximity function) or in situations in which an external current is passed through such a junction.

The work of the authors described above was supported by the Office of Naval Research and an award from the New York State Institute on Superconductivity in conjunction with the New York State Energy Research and Development Authority.

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